Physics Potential of the SPL Super Beam

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Invited talk at the Nufact02 Workshop, Imperial College of Science, Technology and Medicine, London, July 2002.

Abstract. Performances of a neutrino beam generated by the CERN SPL proton driver are computed considering a 440 kton water Čerenkov detector at 130 km from the target. θ_{13} sensitivity down to 1.2° and a δ sensitivity comparable to a Neutrino Factory, for $\theta_{13} \geq 3^{\circ}$, are within the reach of such a project.

1. Introduction

The planned Super Proton Linac (SPL) is a 2.2 GeV proton beam of 4 MW power [1] working with a repetition rate of 75 Hz delivering $1.5 \cdot 10^{14}$ protons per pulse (10^{23} protons on target (pot) in a 10^7 s conventional year). It could be the first stage of a CERN based Neutrino Factory or of a Beta Beam.

Studies of the capabilities of a neutrino beam generated by SPL have been already published in [2], [3], in this paper the fluxes and the overall physical performances will be reviewed in the light of the new design of the beam optics, specially optimized for the SuperBeam needs [4]. They are computed for a gigantic water Čerenkov detector, as the proposed UNO detector [5] (440 kton fiducial) located in the Modane laboratory under the Frejus tunnel at a baseline of 130 km from CERN.

The SPL SuperBeam capabilities would constitute a natural follow-up of the JHF phase I experiment [6], with excellent sensitivity on θ_{13} (section 3) and good sensitivity on the CP phase δ (section 4). Furthermore the SPL SuperBeam could be used to complement the results of a Neutrino Factory experiment, helping in resolving the ambiguities, as discussed in [7]; or could be combined with a Beta Beam, as discussed in [8].

Signal efficiency and backgrounds have already been discussed in [2], they are computed by using the NUANCE neutrino event generator [9] and reconstructing the events with standard SuperKamiokande algorithms, with the addition of improved π° rejection algorithms. They can be summarized as signal efficiency $\epsilon \simeq 70\%$ and π° and μ/e background rejection, normalized to the non oscillated ν_{μ} charged current interactions, $f_B^{\pi^{\circ}} = 4.2 \cdot 10^{-4}$, $f_B^{\mu/e} = 3 \cdot 10^{-3}$.

2. Fluxes

Details of the new beam optics can be found in [4]. The use of an horn and a reflector increases by $\sim 40\%$ the overall ν_{μ} flux with respect to the former single horn optics, slightly increases the ν_{e} contamination, while the $\overline{\nu}_{\mu}$ and $\overline{\nu}_{e}$ contaminations are reduced by $\sim 30\%$.

The length of the decay tunnel has been re-optimized having in mind CP searches more than θ_{13} . Table 1 reports details of the beam properties as function of the length of the decay tunnel, including the sensitivity on θ_{13} for a 2200 kton/year exposure. In spite of the fact that the θ_{13} sensitivity is maximum for the lowest length (20 m), a 60 m decay length is preferred because of the lower $\overline{\nu}_{\mu}$ contamination, that results in a better CP sensitivity. The neutrino spectra for the π^+ and π^- focussed beams are displayed in Fig. 1.

	$\pi^+ focus$			$\pi^- focus$			$\pi^+ focus$
Length	$ u_{\mu}$	ν_e	$\overline{ u}_{\mu}$	$\overline{ u}_{\mu}$	$\overline{ u}_e$	ν_{μ}	θ_{13}
(m)	$(\nu/m^2/yr)$	(%)	(%)	$(\nu/m^2/yr)$	(%)	(%)	(90%CL)
	(@50 km)			(@50 km)			(2200 kton/yr)
20	$2.43 \cdot 10^{+12}$	0.38	1.71	$1.73 \cdot 10^{+12}$	0.41	3.9	1.20
60	$3.23 \cdot 10^{+12}$	0.67	1.50	$2.25 \cdot 10^{+12}$	0.70	3.3	1.25
100	$3.35 \cdot 10^{+12}$	0.76	1.62	$2.33 \cdot 10^{+12}$	0.79	3.3	1.30
20 (old)	$1.71 \cdot 10^{+12}$	0.36	2.4	$1.12 \cdot 10^{+12}$	0.38	5.6	1.47

Table 1. Neutrino fluxes and contamination for different values of the decay tunnel length. The last line refers of the single horn optics of ref. [2].

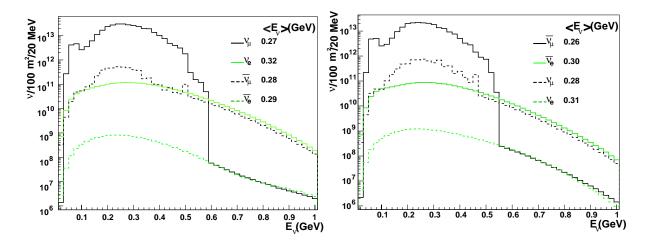


Figure 1. Neutrino spectra for the π^+ (left) and the π^- (right) focussed beam for a decay tunnel length of 60 m.

3. Sensitivity on θ_{13}

The θ_{13} sensitivity is computed assuming $\delta = 0$, solar SMA solution, $\delta m_{23}^2 = 2.5 \cdot 10^{-3}$ eV², $\theta_{23} = 45^{\circ}$ and 5 years of data taking. These are the standard benchmark assumptions used by similar projects [6], [10].

Fig. 2 shows the θ_{13} sensitivity (90% CL) in case of no signal and summarizes the event rate computed for $\theta_{13} = 2^{\circ}$. The experiment would have sensitivity down to $\theta_{13} = 1.2^{\circ} (\sin^2 2\theta_{13} = 1.75 \cdot 10^{-3})$

4. CP sensitivity

CP sensitivity is computed assuming a 2 year run with the π^+ focussed beam and 8 years with the π^- focussed beam. This sharing is motivated by the unfavorable cross section ratio $\overline{\nu}_e / \nu_e \sim 1/6$ at 300 MeV.

A 10% error on the solar δm^2 and $\sin^2 2\theta$, as expected from the KamLAND experiment [11] and a 2% error on the atmospheric δm^2 and $\sin^2 2\theta$, as expected from the JHF neutrino experiment [6] are taken into account. Correlations between θ_{13} and

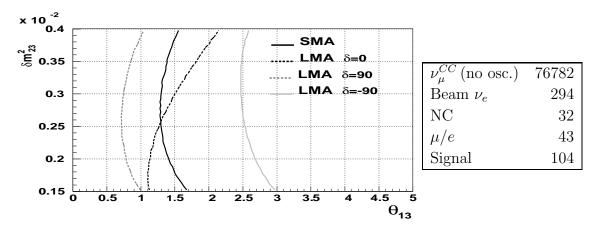


Figure 2. Left: θ_{13} sensitivity (90%CL) computed for a 2200 kton exposure, under the solar SMA solution or under LMA and different δ values. Right: Number of events for the same exposure, SMA solution, in case of $\theta_{13} = 2^{\circ}$.

 δ are fully accounted for, while the sign δm_{13}^2 and the $\theta_{23}/(\pi/2 - \theta_{23})$ ambiguities are not considered. A systematic error of 2% is accounted for the signal efficiency and background normalization, as discussed in [2].

Solutions for different values of δ and θ_{13} , Fig.3-left, show very small correlations between the two parameters.

Since the sensitivity to CP violation heavily depends on the true value of δm_{12}^2 and θ_{13} , we prefer to express the CP sensitivity for a fixed value of δ in the δm_{12}^2 , θ_{13} parameter space. The CP sensitivity to separate $\delta = 90^\circ$ from $\delta = 0^\circ$ at the 99%CL as a function of δm_{12}^2 and θ_{13} , following the convention of [12], is plotted in Fig. 3-right.

It is fair to say that SPL SuperBeam CP sensitivity approaches the Neutrino Factory sensitivity in the parameter space that will be explored by the JHF experiment: $\theta_{13} \geq 2.3^{\circ}$.

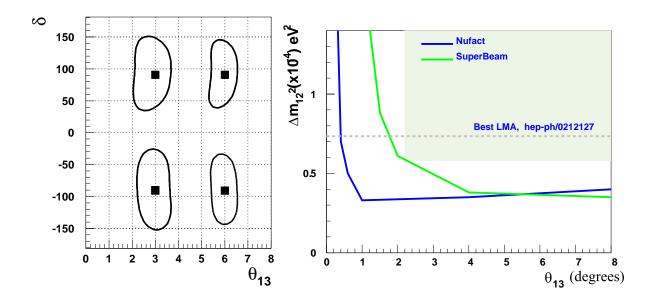


Figure 3. Left: $\theta_{13} - \delta$ fits (99% CL) computed for $\delta m_{12}^2 = 10^{-4} \text{ eV}^2$, $\sin^2 2\theta_{12} = 0.8$. The squares indicate the starting points. Right: CP sensitivity of the SPL-SuperBeam, see text, compared with a 50 GeV Neutrino Factory producing $2 \cdot 10^{20} \mu$ decays/straight section/year, and two 40 kton detectors at 3000 and 7000 km [12]; the shaded region corresponds to the allowed LMA solution and the θ_{13} sensitivity of JHF.

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